

The closure of a shallow tidal inlet promoted by infragravity waves

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Key Points:

- Under storm waves, infragravity waves up to 0.4 m-high propagate through a shallow inlet at flood but are blocked at ebb.
- At the passage of an infragravity wave crest, currents peak over 2.5 m.s^{-1} and sand fluxes increase by two orders of magnitude.
- Over a tidal cycle, large amounts of sand accumulate at the lagoon entrance, which damps tidal propagation and promotes inlet closure.

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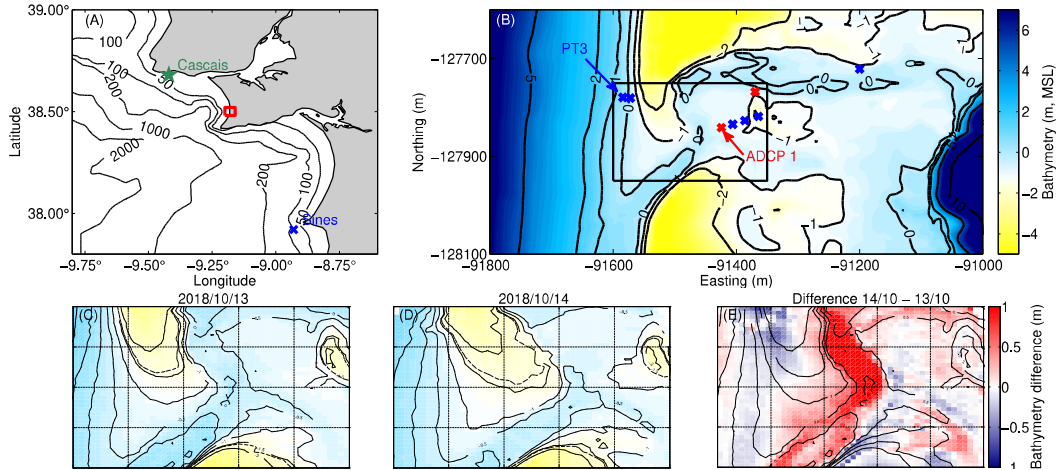
Abstract

Tidal inlets connect the ocean to inner water bodies and are present worldwide. Shallow inlets display fast morphological changes, due to complex interactions between tides, waves and shallow depths. Their closure is commonly observed under storm waves, but the underlying processes remain only partly understood. Here, we present new field evidence that infragravity waves contribute to the closure of shallow inlets. The analysis of new field data collected at a shallow inlet under storm waves reveals that infragravity waves up to 0.4 m-high can propagate inside the lagoon during flood, but are blocked by opposing currents during ebb. At the passage of an infragravity wave crest, currents peak over 2.5 m.s^{-1} and increase instantaneous sand fluxes by two orders of magnitude. Large accumulations of sand at the lagoon entrance damp tidal propagation until full inlet closure. This mechanism provides a new explanation for the closure of shallow inlets observed worldwide.

1 Introduction

Tidal inlets connect the ocean to inner water bodies and concentrate socio-economic and environmental challenges, such as the maintenance of navigation routes and water quality, the mitigation of flooding risk in the backbarrier lagoons and the stabilization of often urbanized adjacent shorelines. Small tidal inlets are particularly dynamic coastal systems, due to the combination of tides, waves and the presence of shallow channels and sandbanks. These dynamics are still only partly understood and drive fast and large morphological changes, making tidal inlet behaviour hard to predict. Tidal inlets will also be substantially impacted by ongoing climate changes [Duong *et al.*, 2016], so that their sustainable management constitutes a major challenge worldwide. Among the fast morphological changes that these coastal systems exhibit, the closure of intermittent inlets is commonly observed in many regions of the world [McSweeney *et al.*, 2017], such as Southwest Europe [Moreno *et al.*, 2010; Fortunato *et al.*, 2014; González-Villanueva *et al.*, 2017], North America [Hanes *et al.*, 2011; Orescanin and Scooler, 2018], West [Anthony *et al.*, 2002] and South Africa [Cooper, 2001] and Australia [Ranasinghe and Pattiaratchi, 2003]. Yet, the underlying processes are still not totally understood while process-based numerical models usually fail to reproduce this behaviour [e.g. Bertin *et al.*, 2009].

Williams and Stacey [2016] performed field measurements at a shallow inlet in Northern California and identified fluctuations in water levels and current velocities with periods of the order of the minute. These fluctuations were interpreted as infragravity waves (hereafter referred to as IG waves), which correspond to long waves associated with the presence of groups in incident short waves [see Bertin *et al.*, 2018, for a recent review]. These authors also showed that velocities associated with IG waves were of the same order of magnitude or even larger than tidal currents, although they rapidly decreased after the beginning of the ebb. Bertin and Olabarrieta [2016] combined field observations and numerical modelling at the Albufeira Lagoon Inlet (SW Portugal, Fig. 1-A) and revealed the occurrence of similar oscillations in water levels and current velocities in the IG frequency band. While these fluctuations were present over the ebb-tidal delta during the whole tidal cycle, they only appeared between the beginning of the flood and up to two hours after high tide inside the lagoon. This behavior was explained by IG wave blocking due to strong counter tidal currents in shallow water depths. Bertin and Olabarrieta [2016] proposed that the occurrence of these fluctuations mostly during flood and not during ebb could promote flood dominance in the lagoon and possibly contribute to their closure during storms. The present study aims at verifying this hypothesis, based on an unpublished field dataset collected at the Albufeira Lagoon Inlet under storm waves.



65 **Figure 1.** (A) Locations of the study area in Portugal (red box) and the nearest operating wave
 66 buoy (blue cross) and tide gauge (green star) during the field campaign. (B) Detailed bathymetry
 67 of the study area (m relative to mean sea level, hereafter MSL), showing the location of the ADCPs
 68 mounted with OBSs (red crosses) and the pressure transducers (blue crosses). Topo-bathymetry of
 69 the inlet region surveyed on the 13/10/2018 (C), the 14/10/2018 (D) and difference between both
 70 (E). Coordinates of (B)-(E) are in meters (coordinates are in ETRS89-PTTM06) and the thick
 71 dotted line in (C)-(D) corresponds to the maximum sustained water level reached inside the lagoon
 72 during the field campaign ($z = -1.6$ m).

73 2 Field campaign and data analysis

74 The study area corresponds to the Albufeira Lagoon, a 1.3 km^2 lagoon located along
 75 the West coast of Portugal and connected to the sea through a shallow, wave-dominated and
 76 intermittently closed inlet [Fortunato *et al.*, 2014]. Tides range from 0.55 to 3.86 m seaward
 77 of the inlet but are strongly distorted and attenuated inside the lagoon. The offshore annual-
 78 mean significant wave height is about 1.9 m but can exceed 10.0 m during storms [Fortunato
 79 *et al.*, 2017]. The sediments in the region of the inlet are coarse, with a median grain size
 80 d_{50} ranging from 0.55 to 0.75 mm [Fortunato *et al.*, 2014].

81 A field campaign was carried out in the Albufeira Lagoon Inlet between the 12/10/2018
 82 and the 16/10/2018 under storm waves associated with the ex-tropical hurricane Leslie,
 83 which made landfall about 150 km to the North of the study area. According to the nearest
 84 operating wave buoy during the campaign (Sines, Fig. 1-A), the offshore spectral significant
 85 wave height $H_{m0,G}$ reached a maximum of 5 m and the peak wave period T_p increased from
 86 12 s to 16 s (Fig. 2-B). Six pressure transducers (hereafter PT) continuously measuring at
 87 4 Hz and buried about 0.05 m in the bed were deployed on both sides of the inlet and two
 88 high-resolution acoustic Doppler current profilers (hereafter ADCP) were deployed around
 89 the flood delta (Fig. 1-B). These two ADCPs were mounted with PTs and optical backscatter
 90 sensors (hereafter OBS) measuring at 2 Hz at the location of the first measuring cell (the
 91 blanking distance was set to 0.10 m). At ADCP1, the first measuring cell was located 0.25
 92 m from the bed (that is 1.16 m above mean sea-level, hereafter MSL), which implies that
 93 currents were not measured about 2 hours before and after low tide in the lagoon.

94 Due to weather forecasts suggesting that offshore $H_{m0,G}$ could reach 10 m, we expected
 95 dramatic morphological changes near the inlet channel during the peak of the storm and so
 96 we did not run the risk of deploying instruments in this area. Each sensor was positioned with

97 a differential GNSS, which provides horizontal and vertical precisions of a few centimeters.
 98 This GNSS was also used to survey the intertidal topography of the inlet region at each low
 99 tide on a daily basis.

100 For each sensor, bottom pressure measurements were first corrected for sea level atmo-
 101 spheric pressure measured at the nearby meteorological station of Cascais (Fig. 1-A). The
 102 entire record was split into consecutive bursts of 20 minutes and only the bursts in which
 103 the sensor was continuously submerged were considered. Bottom pressure energy density
 104 spectra $E_p(f)$ were computed using Fast Fourier Transforms, with 10 Hanning-windowed
 105 segments, which results in 20 degrees of freedom with a frequency resolution of 0.0046 Hz.
 106 These pressure spectra were then converted into elevation spectra $E(f)$ considering linear
 107 wave theory. According to a sensitivity analysis, accounting for tidal currents in the disper-
 108 sion relation has only marginal effects on the correction. The spectral significant wave height
 109 H_{m0} was computed as:

$$H_{m0} = 4\sqrt{m_0} \quad (1)$$

110 with

$$m_0 = \int_{f_{min}}^{f_{max}} E(f)df \quad (2)$$

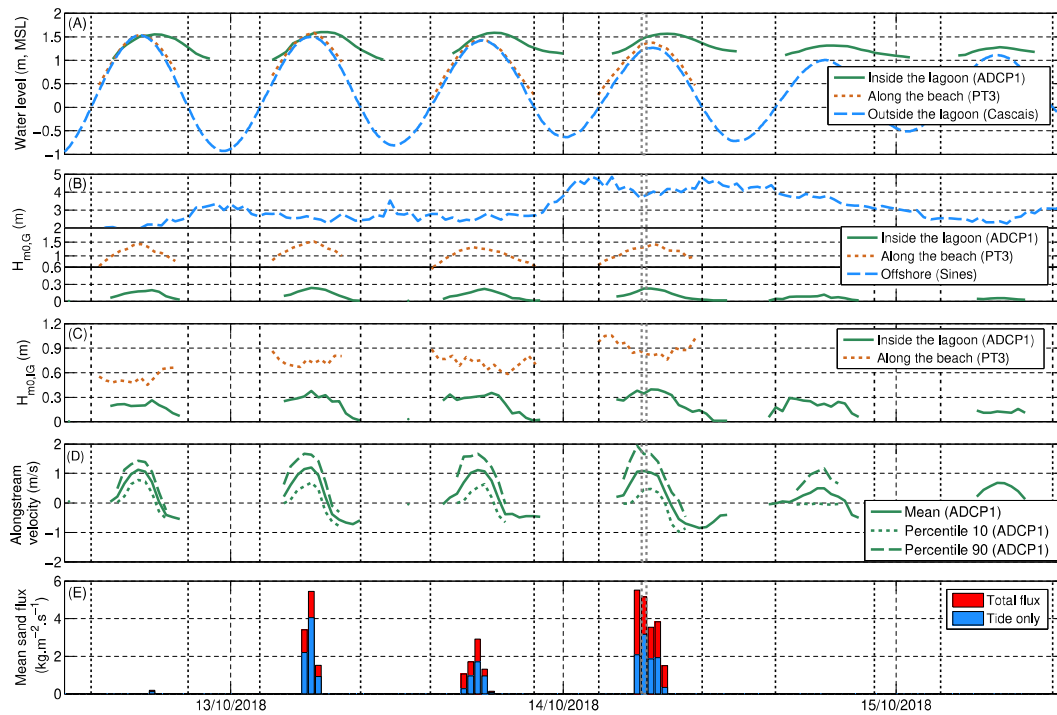
111 In the gravity band, (f_{min}, f_{max}) were set to (0.04, 0.5) Hz. A fixed frequency cutoff
 112 between short and IG waves was considered because incident wave periods varied little
 113 during the field campaign. In the IG band, $H_{m0,IG}$ was integrated between 0.0046 and 0.04
 114 Hz.

115 The ADCPs were set to measure 10 minute-mean velocity profiles with a 0.1 m ver-
 116 tical resolution followed by 20 minutes of continuous measurements at 2 Hz, where current
 117 velocities were averaged over a 0.5 m-thick cell. Depth-integrated current velocities were
 118 computed based on velocity profiles while their short-term variability was investigated us-
 119 ing 20 minute-long samples at 2 Hz. As the OBSs were synchronized to the ADCPs, they
 120 provided data with the same temporal structure. The OBS signal intensity was converted
 121 into sediment concentration after a calibration in the laboratory, using sediment samples
 122 collected at the location of each ADCP.

123 3 Results

124 Digital elevation models (hereafter DEM) with a spatial resolution $\Delta x = 5$ m were de-
 125 rived from the topographic surveys. The comparison between the 13/10/2018 and 14/10/2018
 126 DEMs reveals dramatic morphological changes (Fig. 1-C to E). First, the two sandspits
 127 bounding the inlet migrated towards the lagoon by 30 to 50 m. The seaward side of these
 128 sandspits suffered substantial erosion, with a lowering of the berm locally reaching 1 m,
 129 which corresponds to a local horizontal retreat over 10 m. Second, the main channel suf-
 130 fered a considerable sediment accretion, exceeding 1 m along the northern sandspit. Over
 131 the surveyed area (black rectangle on Fig. 1-B), 2800 m³ of sediment eroded while 10200 m³
 132 accreted, which results in a clearly positive sediment budget (7400 m³). The morphologi-
 133 cal changes of the inlet after the field campaign were monitored qualitatively until its full
 134 closure based on satellite images originating from the Sentinel-2 mission (Appendix A.1).
 135 The analysis of available images reveals that morphological changes of the lagoon entrance
 136 remained very modest (Fig. A.1), without any clear trend since the large sediment accretion
 137 that occurred during the storm Leslie. On the 01/11/2018, the main channel width was of
 138 the order of 10 m and apparently very shallow while the lagoon was totally closed on the
 139 09/11/2018.

145 In order to better understand the causes for these morphological changes, water levels,
 146 waves and currents measured during the field campaign were further analyzed. Fig. 2-A



140 **Figure 2.** Hydro-sedimentary bulk parameters during the field campaign: (A) Water level
 141 variations around mean sea level (m); (B) significant wave height of short waves (m); (C) significant
 142 wave height of IG waves (m); (D) mean and percentile 10 and 90 of low-pass filtered alongstream
 143 current velocities and (E) mean sand fluxes as observed and owing to tidal currents only. The grey
 144 dotted rectangle corresponds to the time of the time-series presented in Fig. 3.

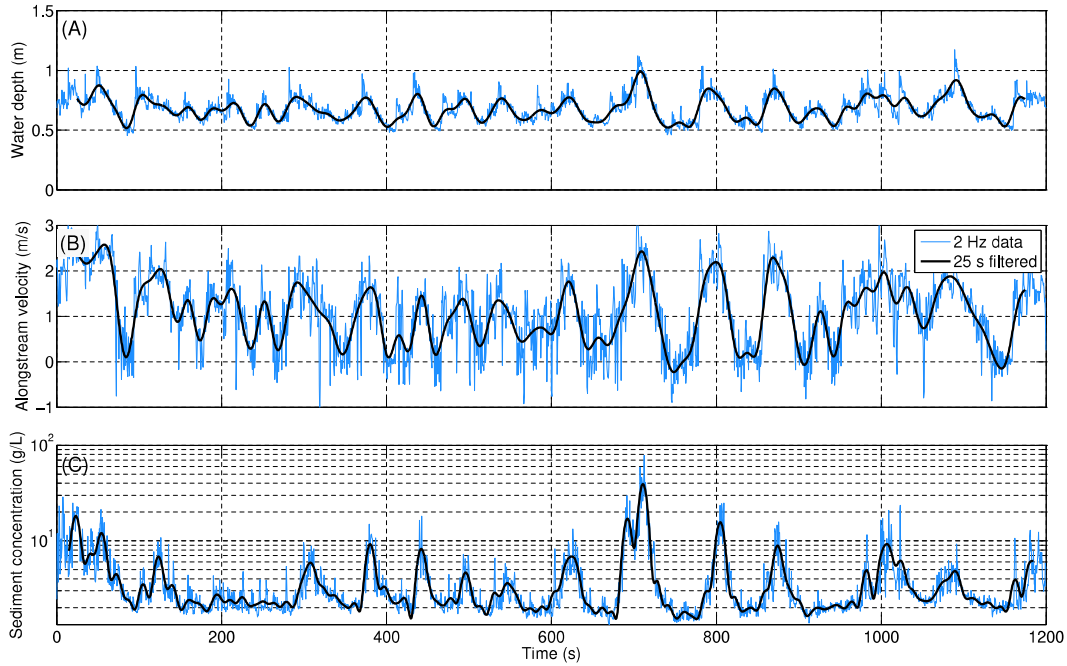
147 shows water levels measured at the nearby tide gauge of Cascais (Fig. 1-A), along the
 148 sandspit outside the lagoon (PT3, Fig. 1-B) and inside the lagoon (ADCP1, Fig. 1-B),
 149 all vertically referenced with respect to MSL. The comparison between these data reveals
 150 firstly that water levels at PT3 were slightly higher than at Cascais (by up to 0.35 m),
 151 which is explained by the development of a wave setup along the coast. Inside the lagoon,
 152 water levels were also higher than at Cascais but the tidal range was dramatically reduced,
 153 a process already described by *Dodet et al.* [2013]. In details, the ratio between the tidal
 154 range inside the lagoon and at Cascais dropped from 0.3 on the 13/10/2018 to 0.15 on the
 155 15/10/2018. Tides inside the lagoon were also strongly distorted, with floods lasting on
 156 average 4h00 and ebbs lasting 8h20.

157 In deep water, $H_{m0,G}$ increased from 2 m at the beginning of the campaign to 5 m at
 158 the peak of the storm (Fig. 2-B) while the peak wave period T_p increased from 11 s to 16 s
 159 (not shown). Along the adjacent beaches, $H_{m0,G}$ was depth-limited and therefore tidally
 160 modulated, with maximum values at high tide ranging from 1.3 m to 1.5 m (Fig. 2-B). Note
 161 that this sensor was always located in the inner surf zone, so that waves were certainly
 162 higher at the breaking point, although we were not able to deploy any sensor there. Inside
 163 the lagoon, $H_{m0,G}$ decreased drastically and ranged from about 0.2 m during the first four
 164 tidal cycles to 0.1 m or less during the last two tidal cycles. $H_{m0,G}$ increased with the water
 165 depth, suggesting that short waves were also depth-limited inside the lagoon. However, the
 166 ratio between $H_{m0,G}$ and the local water depth ranges from 0.15 to 0.40, which suggests that
 167 depth-limitation was not local but rather occurred over the ebb-delta and the southern part
 168 of the northern sandspit. Also, the time series of individual short wave height were positively
 169 correlated with water depth fluctuations associated with IG waves (mean $R^2 = 0.47 \pm 0.12$,
 170 where \pm corresponds to one standard deviation), which is explained by the fact that short
 171 waves were less depth-limited when traveling over the crest of an IG wave. Finally, for a
 172 given water level, short waves were substantially higher during flood than ebb, a behaviour
 173 that was already observed by *Dodet et al.* [2013] and *Bertin and Olabarrieta* [2016] and
 174 explained by wave blocking due to strong opposing ebb currents in shallow depth.

175 Along the adjacent beaches, $H_{m0,IG}$ increased from 0.6 m at the beginning of the
 176 campaign to 1.1 m at the peak of the storm (Fig. 2-C). $H_{m0,IG}$ were slightly tidally modu-
 177 lated, but unlike short waves, they were maximum close to the waterline. Inside the lagoon,
 178 $H_{m0,IG}$ increased from 0.2 m during the first tidal cycle to 0.4 m at the peak of the storm.
 179 Similar to short waves, IG waves almost disappeared after mid-ebb, which was already
 180 observed at this site by *Bertin and Olabarrieta* [2016] and explained by strong opposing
 181 currents in shallow depth.

182 At the location of ADCP1 (Fig. 1-A), 20 minute-mean current velocities reached
 183 1.1 m.s^{-1} during flood and -0.8 m.s^{-1} during ebb for the first four tidal cycles. During the
 184 last two tidal cycles, these values dropped to 0.5 m.s^{-1} and -0.5 m.s^{-1} , respectively (Fig.
 185 2-D). The shorter-term variability of alongstream currents was characterized based on the
 186 mean current speed U_m and their 10 and 90 percentiles (Fig. 2-D). These percentiles show
 187 that, during the first three tidal cycles, maximum current speed values exceeded U_m by 50%
 188 about 10% of the time while minimum values were 50% lower than U_m about 10% of the
 189 time. Around the peak of the storm, maximum values exceeded U_m by almost 100% about
 190 10% of the time while currents were nearly canceled 10% of the time.

195 A time-series of instantaneous water levels and alongstream current velocities around
 196 the peak of the storm is shown on Figs. 3-A and -B. The low-pass filtered water levels
 197 reveal the presence of IG waves of individual height ranging from 0.1 m to 0.5 m. The crests
 198 of these IG waves are associated with peaks in alongstream velocities that reach 2.5 m.s^{-1}
 199 while the passage of their troughs coincides with periods where flood currents are strongly
 200 reduced or even briefly reverted (e.g., see Fig. 3-C around $t = 750 \text{ s}$). Low-pass filtered
 201 water level variations and alongstream velocities are highly correlated, with a correlation
 202 coefficient $r = 0.88$ ($p = 0$). Sediment concentrations measured at the bottom of the first
 203 ADCP cell also display large fluctuations and range from 0.5 g.l^{-1} to over 40 g.l^{-1} . In details,



191 **Figure 3.** Selected time-series at the ADCP1 starting at 6h00 on the 14/10/2018: (A) Water
 192 depth (m); (B) alongstream current velocities and (C) sediment concentration. The blue thin line
 193 corresponds to the raw data at 2 Hz and the black thick line to low-pass filtered data with a 0.04 Hz
 194 frequency cutoff.

204 large sediment concentrations systematically match IG wave crests and associated peaks in
 205 alongstream currents. Low-pass filtered suspended sediment concentrations (hereafter SSC)
 206 and alongstream velocities are also well correlated, with a correlation coefficient $r = 0.61$
 207 ($p = 0$). According to *van Rijn* [2007], the transport of sediment particles in suspension is
 208 proportional to $U^{3.4}$, where U is the current velocity, which implies that SSC is proportional
 209 to $U^{2.4}$. Considering $U^{2.4}$ instead of U yields higher correlations between alongstream
 210 currents and SSC, with a correlation coefficient reaching $r = 0.76$ ($p = 0$). In order to
 211 quantify the contribution of the velocity peaks driven by IG waves to suspended sediment
 212 transport, we compared suspension sand fluxes as observed and suspension sand fluxes that
 213 would be driven by tidal currents only (i.e. using U_m), assuming a $U^{3.4}$ dependency (Fig.
 214 2-E). This analysis confirms that peaks in current velocity driven by IG waves strongly
 215 enhance suspended sand fluxes compared to a situation with tidal currents only, from 35%
 216 to 63% during the first two tidal cycles to more than 400% around the storm peak (210%
 217 on average over the field campaign), when IG waves were the largest. The contribution of
 218 short waves to suspended sediment transport was also evaluated, considering the transport
 219 formula of *van Rijn* [2007] fed with low-pass filtered time series of water levels, current
 220 velocities and individual short wave heights. The comparison between a situation with and
 221 without short waves suggests that short waves increase suspended sand transport by 21%
 222 ($26 \pm 7\%$ on average over the field campaign). Such a modest contribution is related to the
 223 strong decrease of short waves inside the lagoon, with $H_{m0,G}$ ranging from 0.06 to 0.24 m
 224 at ADCP1 over the field campaign.

225 4 Discussion and conclusions

226 *Bertin et al.* [2009] showed that the dynamics of such shallow inlets was controlled by
 227 the balance between tidal asymmetry, which promotes inlet enlargement, and wave-driven

228 processes, which promote sediment accretion at the lagoon entrance. *Bertin and Olabarrieta*
 229 [2016] suggested that IG waves could also have a key contribution to inlet dynamics and
 230 closure. The present dataset provides a unique opportunity to verify this hypothesis.

231 The analysis of water levels and current velocities measured at the Albufeira Lagoon
 232 showed that the passage of IG waves was associated with peaks in current velocities, tem-
 233 porarily exceeding 2.5 m.s^{-1} . These peaks increased instantaneous sand fluxes by up to
 234 two orders of magnitude. Once integrated over 20 minute-long samples, measured sand
 235 fluxes were estimated to be on average 200% larger than in a situation without IG waves,
 236 and as much as 400% when IG waves were largest. This phenomenon only occurs during
 237 flood, because IG waves were blocked during most of the ebb by opposing currents. We
 238 propose that, once integrated over a tidal cycle, this process directly explains the consider-
 239 able sediment accretion that took place at the entrance of the lagoon during the ex-tropical
 240 hurricane Leslie. Note that, due to the low elevation of the northern sandspit, a large part
 241 of its surface-area was continuously submerged and exposed to flood currents around high
 242 tide (Fig. 1-C). Therefore, the mechanism described above was certainly not limited to the
 243 inlet main channel but was probably active over the sandspit area located below high tide
 244 water level, that is $z = -1.6 \text{ m}$. Additionally, both visual field observations and repetitive
 245 topographic surveys revealed that the northern sandpit suffered morphological changes up
 246 to $z = -2.5 \text{ m}$. Considering *Stockdon et al.* [2006] using a mean surf zone slope $\beta = 0.05$ and
 247 deep water wave conditions as measured in Sines (Fig. 1-A), the runup 2% exceedence $R_{2\%}$
 248 would be of the order of 1.4 m, which yields a maximum runup level around 3.0 m above
 249 MSL. Therefore, overwashes driven by IG waves could explain the large sediment deposition
 250 observed inside the lagoon along northernmost part of the North spit (Fig. 1-E).

251 The topographic surveys carried out after the peak of the storm (not shown) as well
 252 as the time series of Sentinel-2 images (A.1) suggest that the lagoon entrance displayed
 253 very limited morphological changes during the two weeks that followed the storm. Water
 254 levels (Fig. 2-A) and current velocities (Fig. 2-D) measured inside the lagoon revealed
 255 that the tidal range and currents decreased by a factor of two after the passage of Leslie.
 256 This combination of topographic and hydrodynamic measurements suggests that the large
 257 sediment accretion at the entrance of the lagoon damped tidal propagation inside the lagoon,
 258 so that the inlet became almost inactive during the two weeks following Leslie, until the
 259 main channel fully closed. We propose that the final closure of the inlet channel could
 260 result from the combination of longshore transport and swash processes [e.g. *Baldock et al.*,
 261 2008], which promoted the development of the steep berm that is usually present along the
 262 beachface.

263 In conclusion, this study demonstrates that the propagation of IG waves in a shallow
 264 tidal inlet can have a fundamental contribution to its hydro-sedimentary dynamics and
 265 largely contributed to its closure. Although based on a single case study, this unreported
 266 mechanism is likely to be valid for most shallow inlets exposed to the open ocean, since IG
 267 waves develop as soon as groups are present in the incident short waves reaching the coast.
 268 Global IG wave climate computed by *Ardhuin et al.* [2014] would suggest that IG waves
 269 should be particularly relevant for the dynamics of shallow inlets located in northwestern and
 270 southwestern America, western Europe, northwestern and southwestern Africa and western
 271 and southern Australia. The relevance of IG waves will also have to be verified at other types
 272 of inlets, including larger and deeper systems, for which the impact of IG waves is totally
 273 unknown according to our knowledge. Up to now, IG waves were usually not represented in
 274 numerical models when applied to tidal inlets: this study suggests that including IG waves
 275 in future modelling studies should improve model predictive skills, and possibly allow to
 276 reproduce for the first time tidal inlet closure without unrealistic parameterizations. The
 277 contribution of swash processes to the final closure of inlet channels will also have to be
 278 analyzed in the future.

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 287 (<http://doi.org/10.5281/zenodo.3141501>), which will be available after an embargo period
 288 corresponding to the completion of this PhD thesis (September 2020). Offshore wave data
 289 in Sines can be obtained through a formal request to the Portuguese Hydrographic Institute.
 290 Water level data in Cascais was provided by the Direção Geral do Território, and can be
 291 obtained at <ftp.dgterritorio.pt/Maregrafos/Cascais>. Sentinel-2 images are made available by
 292 the European Space Agency through the Copernicus Portal. Ana Silva provided a valuable
 293 help on the field at the beginning of the campaign and Thibault Coulombier performed the
 294 calibration of the OBSs in the lab.

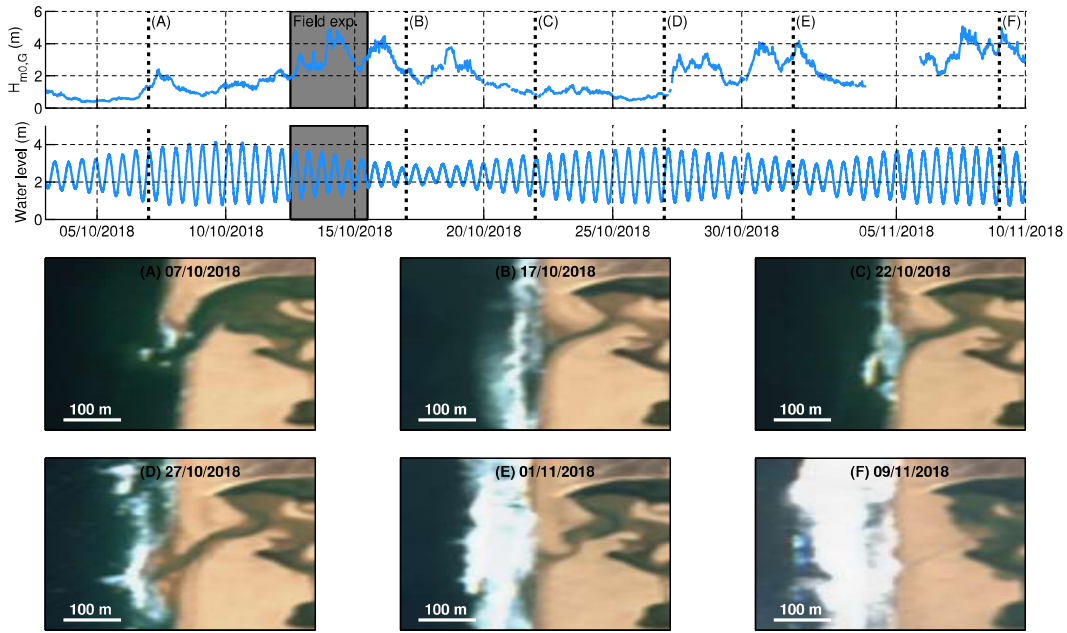
295 **References**

- 296 Anthony, E., L. Oyédé, and J. Lang (2002), Sedimentation in a fluvially infilling, barrier-
 297 bound estuary on a wave-dominated, microtidal coast: The Ouémé River Estuary, Benin,
 298 West Africa, *Sedimentology*, *49*(5), 1095–1112, doi:10.1046/j.1365-3091.2002.00491.x.
- 299 Arduin, F., A. Rawat, and J. Aucan (2014), A numerical model for free infragravity waves:
 300 Definition and validation at regional and global scales, *Ocean Modelling*, *77*, 20–32.
- 301 Baldock, T., F. Weir, and M. Hughes (2008), Morphodynamic evolution of a
 302 coastal lagoon entrance during swash overwash, *Geomorphology*, *95*(3), 398–411, doi:
 303 <https://doi.org/10.1016/j.geomorph.2007.07.001>.
- 304 Bertin, X., and M. Olabarrieta (2016), Relevance of infragravity waves at a wave-dominated
 305 inlet, *Journal of Geophysical Research*, (212), 5418–5435, doi:10.1002/2015JC011444.
- 306 Bertin, X., A. B. Fortunato, and A. Oliveira (2009), A modeling-based analysis of pro-
 307 cesses driving wave-dominated inlets, *Continental Shelf Research*, *29*(5-6), 819–834, doi:
 308 10.1016/j.csr.2008.12.019.
- 309 Bertin, X., A. de Bakker, A. van Dongeren, G. Coco, G. André, F. Arduin, P. Bonneton,
 310 F. Bouchette, B. Castelle, W. Crawford, M. Davidson, M. Deen, G. Dodet, T. Guérin,
 311 K. Inch, F. Leckler, R. McCall, H. Muller, M. Olabarrieta, D. Roelvink, G. Ruessink,
 312 D. Sous, E. Stutzmann, and M. Tissier (2018), Infragravity waves: From driving mecha-
 313 nisms to impacts, *Earth-Science Reviews*, *177*, 774–799.
- 314 Cooper, J. (2001), Geomorphological variability among microtidal estuaries from the wave-
 315 dominated south african coast, *Geomorphology*, *40*(1-2), 99–122, doi:10.1016/S0169-
 316 555X(01)00039-3.
- 317 Dodet, G., X. Bertin, N. Bruneau, A. B. Fortunato, A. Nahon, and A. Roland (2013), Wave-
 318 current interactions in a wave-dominated tidal inlet, *Journal of Geophysical Research:*
 319 *Oceans*, *118*(3), 1587–1605, doi:10.1002/jgrc.20146.
- 320 Drusch, M., U. Del Bello, S. Carlier, O. Colin, V. Fernandez, F. Gascon, B. Hoersch, C. Isola,
 321 P. Laberinti, P. Martimort, A. Meygret, F. Spoto, O. Sy, F. Marchese, and P. Bargellini
 322 (2012), Sentinel-2: ESA’s optical high-resolution mission for GMES operational services,
 323 *Remote Sensing of Environment*, *120*, 25–36, doi:10.1016/j.rse.2011.11.026.
- 324 Duong, T., R. Ranasinghe, D. Walstra, and D. Roelvink (2016), Assessing climate change
 325 impacts on the stability of small tidal inlet systems: why and how?, *Earth-Science Re-*
 326 *views*, *154*, 369–380.
- 327 Fortunato, A., P. Freire, X. Bertin, M. Rodrigues, J. Ferreira, and M. Liberato (2017), A
 328 numerical study of the february 15, 1941 storm in the tagus estuary, *Continental Shelf*
 329 *Research*, *144*, 50–64.

- 330 Fortunato, A. B., A. Nahon, G. Dodet, R. A. Pires, C. M. Freitas, N. Bruneau, A. Azevedo,
331 X. Bertin, P. Benevides, C. Andrade, and A. Oliveira (2014), Morphological evolution of
332 an ephemeral tidal inlet from opening to closure: the Albufeira Inlet, Portugal, *Conti-*
333 *nerental Shelf Research*, *73*, 49–63.
- 334 González-Villanueva, R., M. Pérez-Arlucea, and S. Costas (2017), Lagoon water-level oscil-
335 lations driven by rainfall and wave climate, *Coastal Engineering*, *130*, 34–45.
- 336 Hanes, D., K. Ward, and L. Erikson (2011), Waves and tides responsible for the intermit-
337 tent closure of the entrance of a small, sheltered tidal wetland at San Francisco, CA,
338 *Continental Shelf Research*, *31*(16), 1682–1687.
- 339 McSweeney, S., D. Kennedy, I. Rutherford, and J. Stout (2017), Intermittently closed/open
340 lakes and lagoons: Their global distribution and boundary conditions, *Geomorphology*,
341 *292*, 142–152, doi:<https://doi.org/10.1016/j.geomorph.2017.04.022>.
- 342 Moreno, I., A. Ávila, and M. Losada (2010), Morphodynamics of intermittent coastal lagoons
343 in Southern Spain: Zahara de los Atunes, *Geomorphology*, *121*(3-4), 305–316.
- 344 Orescanin, M., and J. Scooler (2018), Observations of episodic breaching and closure at an
345 ephemeral river, *Continental Shelf Research*, *166*, 77–82.
- 346 Ranasinghe, R., and C. Pattiaratchi (2003), The seasonal closure of tidal inlets: Causes and
347 effects, *Coastal Engineering Journal*, *45*(4), 601–627, doi:10.1142/S0578563403000919.
- 348 Stockdon, H. F., R. A. Holman, P. A. Howd, and A. H. Sallenger Jr (2006), Empirical
349 parameterization of setup, swash and runup, *Coastal Engineering*, *53*, 573–588.
- 350 van Rijn, L. (2007), Unified view of sediment transport by currents and waves. ii: Suspended
351 transport, *Journal of Hydraulic Engineering*, *133*(6), 668–689, doi:10.1061/(ASCE)0733-
352 9429(2007)133:6(668).
- 353 Williams, M. E., and M. T. Stacey (2016), Tidally discontinuous ocean forcing in bar-built
354 estuaries: The interaction of tides, infragravity motions, and frictional control, *Journal*
355 *of Geophysical Research*, *121*, 571–585, doi:10.1002/2015JC011166.

356 **A: Morphological evolution of the inlet based on Sentinel images**

357 The morphological changes of the Albufeira Lagoon Inlet outside the field campaign
 358 were qualitatively monitored based on satellite images obtained from the Sentinel-2 mission.
 359 Sentinel-2 is a large swath, high-resolution, multi-spectral satellite mission launched by
 360 the European Space Agency in June 2015 (Sentinel-2A) and complemented with a second
 361 satellite launched in March 2017 (Sentinel-2B, *Drusch et al. [2012]*). These two satellites
 362 are flying along the same orbit but out of phase and allow for a minimum revolution period
 363 of about 6 days. In the visible domain, the spatial resolution is 10 m.



364 **Figure A.1.** Time series of significant wave height at Sines Buoy and water levels at Cascais
 365 tide gauge (Fig. 1-(A)), where the vertical dashed-dotted black lines correspond to the dates of
 366 selected Sentinel images (A) to (F), showing that considerable morphological changes took place
 367 during the field campaign, after which the inlet remained poorly dynamic until full closure.